

Block 20 — Inductance and Energy

Student Group

First Name	Surname	Matrikel Nr.

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Block 20 — Inductance and Energy

20.0 Intro

20.0.1 Learning objectives

After this 90-minute block, you can

- ...

20.0.2 Preparation at Home

Well, again

- read through the present chapter and write down anything you did not understand.
- Also here, there are some clips for more clarification under 'Embedded resources' (check the text above/below, sometimes only part of the clip is interesting).

For checking your understanding please do the following exercises:

- ...

20.0.3 90-minute plan

1. Warm-up (x min):
 1.
2. Core concepts & derivations (x min):
 1. ...
3. Practice (x min): ...
4. Wrap-up (x min): Summary box; common pitfalls checklist.

20.0.4 Conceptual overview

1. ...

20.1 Core content

20.1.1 Self-Induction

Up to now, we investigated the induction of electric voltages and currents based on the change of an external flux $\frac{d\Phi_{\text{ext}}}{dt}$. For the induced current i_{ind} , we found that it counteracts the change of the external flux (Lenz law).

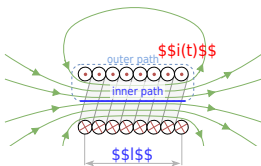
But what happens, when there is no external field - only a coil which creates the flux change itself

(see [figure 1](#))?

Fig. 1: Induction Phenomenons

To understand this, we will investigate the situation for a long coil ([figure 2](#)).

Fig. 2: Self-Induction of a Coil



Given by the [Recap of the fieldline images](#) in Block16, we know that the H -field is given by magnetic voltage $\oint H(t) \cdot dl = N \cdot i$ as:

$$H(t) = \frac{N \cdot i}{l}$$

This leads to the magnetic flux density $B(t)$

$$B(t) = \mu_0 \mu_r \cdot \frac{N \cdot i}{l}$$

Based on the magnetic flux density $B(t)$ it is possible to calculate the flux $\Phi(t)$:

$$\Phi(t) = \int_A \vec{B}(t) \cdot d\vec{A} = \int_A \mu_0 \mu_r \cdot \frac{N \cdot i}{l} \cdot dA = \mu_0 \mu_r \cdot \frac{N \cdot i}{l} \cdot A$$

The changing flux Φ is now creating an induced electric voltage and current, which counteracts the initial change of the current.

This effect is called **Self Induction**. The induced electric voltage u_{ind} is given by:

$$u_{ind} = -N \cdot \frac{d\Phi(t)}{dt} = -N \cdot \frac{d}{dt} \left(\mu_0 \mu_r \cdot \frac{N \cdot i}{l} \cdot A \right) = -N \cdot \mu_0 \mu_r \cdot \frac{N \cdot A}{l} \cdot \frac{di}{dt}$$

$$\boxed{u_{ind} = -\mu_0 \mu_r \cdot N^2 \cdot \frac{A}{l} \cdot \frac{di}{dt}}$$

$$\frac{i}{\frac{d}{dt}} \quad \text{\texttt{\textbackslash\textbackslash \text{for a long coil} \end{align*}}$$

The result means that the induced electric voltage u_{ind} is proportional to the change of the current $\frac{d}{dt}i$.

The proportionality factor is also called **Self-inductance** L (or often simply called inductance).

20.1.2 Inductance

The inductance is another passive basic component of the electric circuit. Besides the ohmic resistor R and the capacitor C , the inductor L is the lump component entailing the inductance.

Generally, the inductance is defined by:
$$L = \left| \frac{u_{\text{ind}}}{\frac{d}{dt}i} \right|$$

The inductance L can also be described differently based on Lenz law $u_{\text{ind}} = - \frac{d}{dt}\Psi(t)$:

$$L = \left| \frac{u_{\text{ind}}}{\frac{d}{dt}i} \right| = \frac{d\Psi(t)/dt}{\frac{d}{dt}i}$$

$$L = \frac{\Psi(t)}{i}$$

One can also consider an inductor a “conservative person”: it does not like to see abrupt changes in the passing current. It reacts to any change in the current with a counteracting voltage since the current change leads to a changing flux and - therefore - an induced voltage. The [figure 3](#) shows an inductor in series with a resistor and a switch (any real switch also behaves as a capacitor, when open). Once the simulation is started, the inductor directly counteracts the current, which is why the current only slowly increases.

The unit of the inductance is $1 \text{ Henry} = 1 \text{ H} = 1 \frac{\text{Vs}}{\text{A}} = 1 \frac{\text{Wb}}{\text{A}}$

Fig. 3: Example of a Circuit with an Inductor

Mathematically the voltages can be described in the following way:

$$u_0 = u_R + u_L \quad \text{\texttt{\textbackslash\textbackslash \text{for a long coil} \end{align*}}$$

20.1.3 Inductance of different Components

Long Coil

In the last sub-chapter, the formula of a long coil was already investigated. By these, the inductance of a long coil is

$$L_{\text{long coil}} = \mu_0 \mu_r \cdot N^2 \cdot \frac{A}{l}$$

Toroidal Coil

The toroidal coil was analyzed in the last chapter (see [magnetic Field Strength Part 1: Toroidal Coil](#)). Here, a rectangular intersection is assumed (see [figure 4](#)).

Fig. 4: Self-Induction of a toroidal Coil

This leads to

$$H(t) = \frac{N \cdot i}{l}$$

with the mean magnetic path length (= length of the average field line) $l = \pi(r_o + r_i)$:

$$H(t) = \frac{N \cdot i}{\pi(r_o + r_i)}$$

The inductance L can be calculated by

$$L_{\text{toroidal coil}} = \frac{\Psi(t)}{i} \quad \&= \quad \frac{N \cdot \Phi(t)}{i}$$

With the magnetic flux density $B(t) = \mu_0 \mu_r H(t) = \mu_0 \mu_r \frac{i \cdot N}{\pi(r_o + r_i)}$ and the cross section $A = \pi(r_o - r_i)$, we get:

$$\begin{aligned} \quad \quad L_{\text{toroidal \; coil}} &= \frac{N \cdot \mu_0 \mu_r \int \mathbf{i} \cdot d\mathbf{l}}{\oint (\mathbf{r}_o + \mathbf{r}_i)} \cdot \frac{h(\mathbf{r}_o) - h(\mathbf{r}_i)}{i} \quad \quad \quad = \frac{N^2 \cdot \mu_0 \mu_r \cdot h(\mathbf{r}_o) - h(\mathbf{r}_i)}{\oint (\mathbf{r}_o + \mathbf{r}_i)} \quad \quad \quad \end{aligned}$$

$$\boxed{L_{\text{toroidal \; coil}}} = \mu_0 \mu_r \cdot N^2 \cdot \frac{h(\mathbf{r}_o) - h(\mathbf{r}_i)}{\oint (\mathbf{r}_o + \mathbf{r}_i)} \quad \quad \quad \end{aligned}$$

20.1.4 Inductances in Circuits

Focus here: uncoupled inductors!

Series Circuits

Based on $L = \frac{\Psi(t)}{i}$ and Kirchhoff's mesh law ($i = \text{const}$) the series circuit of inductions can be interpreted as a single current i which generates multiple linked fluxes Ψ_i . Since the current must stay constant in the series circuit, the following applies for the equivalent inductor of a series connection of single ones:

$$L_{\text{eq}} = \frac{\sum_i \Psi_i}{i} = \sum_i L_i$$

A similar result can be derived from the induced voltage $u_{\text{ind}} = L \frac{di}{dt}$, when taking the situation of a series circuit (i.e. $i_1 = i_2 = i_3 = \dots = i_{\text{eq}}$ and $u_{\text{eq}} = u_1 + u_2 + \dots$):

$$\begin{aligned} u_{\text{eq}} &= u_1 + u_2 + \dots \quad \quad \quad L_{\text{eq}} \frac{di_{\text{eq}}}{dt} \\ &= L_1 \frac{di_1}{dt} + L_2 \frac{di_2}{dt} + \dots \\ &= L_1 \frac{di}{dt} + L_2 \frac{di}{dt} + \dots \quad \quad \quad L_{\text{eq}} \frac{di}{dt} = L_1 \frac{di}{dt} + L_2 \frac{di}{dt} + \dots \quad \quad \quad \end{aligned}$$

Parallel Circuits

For parallel circuits, one can also start with the principles based on Kirchhoff's mesh law:

$$u_{\text{eq}} = u_1 = u_2 = \dots$$

and Kirchhoff's nodal law:

$$i_{\text{eq}} = i_1 + i_2 + \dots$$

Here, the formula for the induced voltage has to be rearranged:

$$\begin{aligned} u_{\text{ind}} &= L \frac{di}{dt} \quad \quad \quad \int u_{\text{ind}} dt = L \cdot i \quad \quad \quad i = \frac{1}{L} \int u_{\text{ind}} dt \quad \quad \quad \end{aligned}$$

By this, we get:

$$\begin{aligned} i_{\text{eq}} &= i_1 + i_2 + \dots = \frac{1}{L_{\text{eq}}} \int u_{\text{eq}} dt \\ &= \frac{1}{L_1} \int u_1 dt + \frac{1}{L_2} \int u_2 dt + \dots \\ &= \frac{1}{L_1} \int u dt + \frac{1}{L_2} \int u dt + \dots = \left(\frac{1}{L_1} + \frac{1}{L_2} + \dots \right) \int u dt \\ &= \frac{1}{L_1} \int u dt + \frac{1}{L_2} \int u dt + \dots \quad \quad \quad \end{aligned}$$

Notice:

The inductor behaves in the parallel and series circuit similar to the resistor.

20.1.5 Energy of the magnetic Field

not covered

20.2 Common pitfalls

- ...

20.3 Exercises

Exercise E1 Self-Induction

(written test, approx. 8 % of a 120-minute written test, SS2024)

2. Determine the inductance of a coil with a radius of 2 cm and 500 turns. The current through the coil changes linearly from 0 A to 3 A in 0.02 ms . The arrangement is located in air ($\mu_{\text{r}} = 1$).
Path

$\mu_0 = 4\pi \cdot 10^{-7} \text{ Vs/Am}$
 $U_{\text{ind}} = 1.32 \text{ mV}$

.. Calculate the (self-)inductance of the coil.

For the linear change of the current the formula of the induced voltage can also be linearized:
$$u_{\text{ind}} = -L \cdot \frac{di}{dt} \quad \Leftrightarrow -L \cdot \frac{\Delta i}{\Delta t} \quad \Leftrightarrow -1.32 \cdot 10^{-3} \text{ V} \cdot \frac{3 \text{ A}}{0.02 \cdot 10^{-3} \text{ s}}$$

The formula for the induction of a long coil is:
$$L = \mu_0 \mu_{\text{r}} \cdot N^2 \cdot \frac{A}{l} = 4\pi \cdot 10^{-7} \text{ Vs/Am} \cdot (500)^2 \cdot \frac{\pi \cdot (2 \cdot 10^{-2} \text{ m})^2}{2 \cdot 10^{-2} \text{ m}}$$

Exercise E11 Self Induction

(written test, approx. 8 % of a 120-minute written test, SS2022)

2. A circuit with a resistor and an inductor is connected to a DC voltage source, which is fused with a circuit breaker.

Sketch the drag torque T_d and disc ω_d if $\omega_m = 0$ with a back EMF $E_b = 63$ V on the induced voltage (induced by) $\omega_m = 3150$ rpm. $\mu_s = 1$ within $1 \sim \mu_s$.
The inner resistance of the motor shall be neglected.)

Path

... Draw the circuit (the circuit breaker can be drawn as a switch), with all voltage and current arrows.

For the maximum voltage on the circuit breaker one has to consider the following:

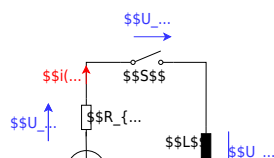
Result

- external voltage of the voltage source U_{ext}
- voltage $U_{\text{ind}}(t)$ induced by the change of the current

The first one is not given in the exercise, and therefore not considered here.

The induced voltage can be calculated by linearizing the following:
$$u_{\text{ind}}(t) = -L \frac{di}{dt} \rightarrow u_{\text{ind}}(t) = -L \frac{\Delta i}{\Delta t}$$

With the given details:
$$u_{\text{ind}}(t) = -L \frac{0 - I}{t_1 - t_0} = 50 \cdot 10^{-6} \frac{63 \text{ A}}{1 \cdot 10^{-6} \text{ s}} = 3150 \frac{\text{Vs}}{\text{A}} \cdot \frac{\text{A}}{\text{s}} = 3150 \frac{\text{V}}{\text{s}}$$



Exercise 4.5.1 Self Induction I

Result $\{\rm r\,m\,r\}=1000\$)$

1. Cylindrical long air coil with $N=390$, winding diameter $d=3.0 \text{ cm}$ and length $l=18 \text{ cm}$

Solution

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Solution: begin{align*} L_1 &= 0.75 \text{ mm} \text{ (height of the liquid)} \end{align*}
```

$$\frac{\hbar^2}{2m} \frac{d^2 \psi}{dx^2} + V(x) \psi = E \psi$$

$$\cdot 10^{-7} \left\{ \frac{H}{m} \right\} \cdot 1 \cdot (390)^2 \cdot \frac{\pi \cdot (0.03 \text{ m})^2}{2} \cdot \frac{1}{0.18 \text{ m}} \end{align*}$$

Exercise 4.5.2 Self Induction II

A cylindrical air coil (length $l=40 \text{ cm}$, radius $r=5.0 \text{ cm}$, and a number of windings $N=300$) passes a current of 30 A . The current shall be reduced linearly in 2.0 ms down to 0.0 A .

What is the amount of the induced voltage u_{ind} ?

$$\begin{align*} |u_{\text{ind}}| &= 33 \text{ V} \end{align*}$$

Solution

The requested induced voltage can be derived by:

$$\begin{align*} L &= \frac{u_{\text{ind}}}{\frac{di}{dt}} \\ \Rightarrow |u_{\text{ind}}| &= L \cdot \left| \frac{di}{dt} \right| \\ &= L \cdot \left| \frac{\Delta i}{\Delta t} \right| \end{align*}$$

Therefore, we just need the inductance L , since $\frac{\Delta i}{\Delta t}$ is defined as 30 A per 2 ms :

$$\begin{align*} L &= \mu_0 \mu_r \cdot N^2 \cdot \frac{A}{l} \\ \end{align*}$$

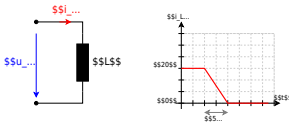
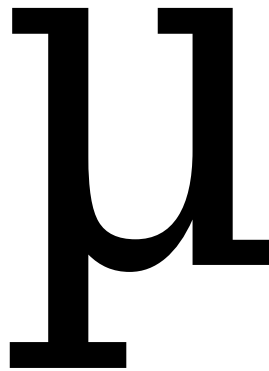
$$\begin{align*} \text{So, the result can be derived as: } |u_{\text{ind}}| &= \mu_0 \mu_r \cdot N^2 \cdot \frac{A}{l} \cdot \left| \frac{\Delta i}{\Delta t} \right| \\ &= 4\pi \cdot 10^{-7} \left\{ \frac{H}{m} \right\} \cdot 1 \cdot (300)^2 \cdot \frac{\pi \cdot (0.05 \text{ m})^2}{0.40 \text{ m}} \cdot \frac{30 \text{ A}}{2 \text{ ms}} \end{align*}$$

Exercise 4.5.3 Self Induction III

A coil with the inductance $L=20 \text{ }\mu\text{H}$ passes a current of 40 A . The current shall be reduced linearly in $5 \text{ }\mu\text{s}$ down to 0 A (see [figure 5](#)).

- What is the amount of the induced voltage u_{ind} ?
- Sketch the course of $u_{\text{ind}}(t)$!

Fig. 5: Circuit and timing Diagram



Embedded resources

Explanation (video): ...

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